

Towards a Solution to the So-Called 3-Generations Problem of Elementary Particles

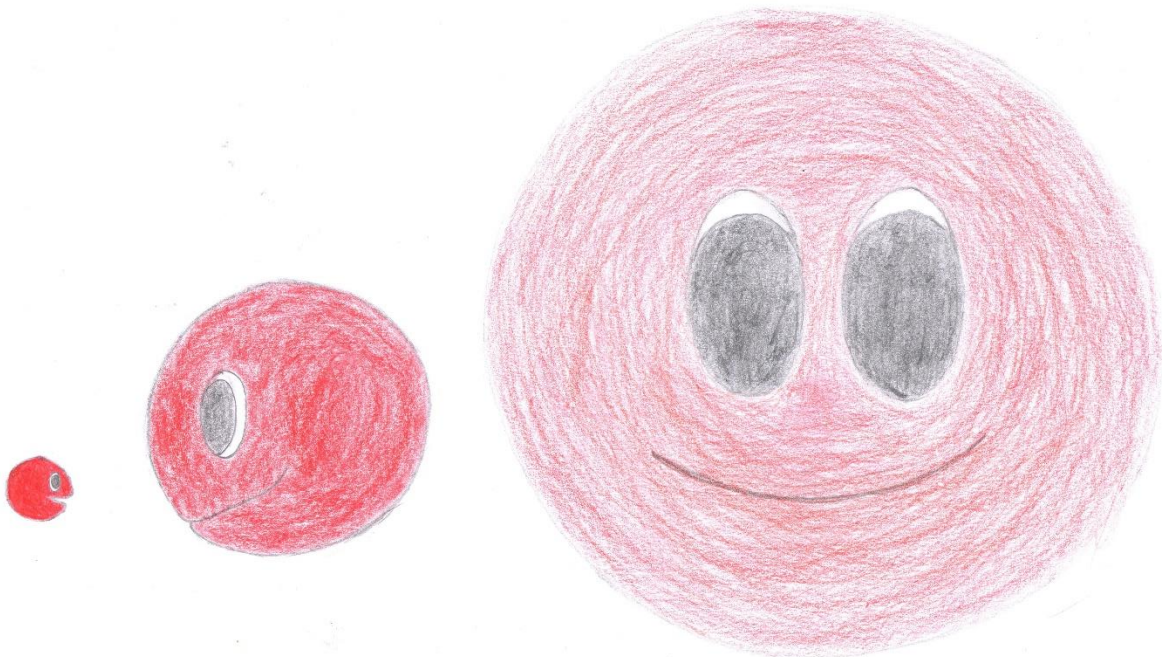
Or: Why are There 3 masses for Charged Leptons, Neutrons and Quarks?

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The Three Generations of Particles, by Livia Schwarzer

Theory: The Ricci Scalar for a Scaled Metric

In [1, 2, 3] (and quite some other papers by this author) it was demonstrated how we can find a direct extraction of the Klein-Gordon, the Schrödinger and the Dirac equation from the Ricci scalar R^* of a modified metric of the kind $G_{\alpha\beta} = F[f] \cdot g_{\alpha\beta}$. Namely: With a yet arbitrary scalar function $F[f]$ the corresponding modified Ricci scalar R^* reads:

$$R^* = \frac{1}{F[f]^3} \cdot \left(\begin{aligned} & \left(C_{N1} \cdot \left(\frac{\partial F[f]}{\partial f} \right)^2 - C_{N2} \cdot F[f] \cdot \frac{\partial^2 F[f]}{\partial f^2} \right) \cdot \overbrace{(\tilde{\nabla}_g f)^2}^{=f_{,\alpha} g^{\alpha\beta} f_{,\beta}} \\ & - C_{N2} \cdot F[f] \cdot \frac{\partial F[f]}{\partial f} \cdot \Delta_g f \end{aligned} \right) \quad (1)$$

$$\xrightarrow{n=4} = \frac{1}{F[f]^3} \cdot \left(\left(\frac{3}{2} \cdot \left(\frac{\partial F[f]}{\partial f} \right)^2 - 3 \cdot F[f] \cdot \frac{\partial^2 F[f]}{\partial f^2} \right) \cdot (\tilde{\nabla}_g f)^2 - 3 \cdot F[f] \cdot \frac{\partial F[f]}{\partial f} \cdot \Delta_g f \right)$$

Thereby the C_{N1} and C_{N2} are constants depending on the number of dimensions n of the space-time being considered. For illustration we also gave the case $n=4$ (n =number of dimensions). Thereby f is to be understood as a function of the coordinates $f=f[x_0, x_1, x_2, \dots]$. It was shown in [1 - 6] (especially [2]) that we can write:

$$C_{N1} = -\frac{3}{2} + \frac{7}{4} \cdot n - \frac{1}{4} \cdot n^2, \quad (2)$$

$$C_{N2} = n - 1$$

Demanding certain conditions for the function $F[f]$ of f then gives us Dirac or Klein-Gordon equations [1, 2, 3].

So, for instance, by setting $F[f]$ to fulfill the following condition:

$$C_{N1} \cdot \left(\frac{\partial F[f]}{\partial f} \right)^2 - C_{N2} \cdot F[f] \cdot \frac{\partial^2 F[f]}{\partial f^2} = 0, \quad (3)$$

we obtain:

$$R^* = -\frac{1}{F[f]^3} \cdot C_{N2} \cdot F[f] \cdot \frac{\partial F[f]}{\partial f} \cdot \Delta_g f$$

$$\xrightarrow{n=4 \ \& \ F[f]=\left(f+\frac{C_f}{M^2}\right)^2} = -\frac{1}{F[f]^2} \cdot 3 \cdot \frac{\partial F[f]}{\partial f} \cdot \Delta_g f \xrightarrow{F[f]=\left(f+\frac{C_f}{M^2}\right)^2} = -\frac{6}{\left(f+\frac{C_f}{M^2}\right)^3} \cdot \Delta_g f \quad (4)$$

$$\Rightarrow \frac{R^*}{6} \cdot \left(f + \frac{C_f}{M^2}\right)^3 + \Delta_g f = 0$$

Towards an Explanation for the 3-Generations Problem (?)

Now we know that the classical Klein-Gordon equation always is of the kind:

$$-M^2 f + \Delta_g f = (-M^2 + \Delta_g) f = 0$$

$$\text{with : } M^2 \equiv \frac{m^2 \cdot c^2}{\hbar^2} \quad (5)$$

(c... speed of light in vacuum, m... rest mass, \hbar ...reduced Planck constant).

Comparing the last line of (4) with the first line of (5) tells us, that finding a linearization for the expression $\frac{R^*}{6} \cdot \left(f + \frac{C_f}{M^2} \right)^3$ with respect to the function f, would immediately give us the classical

Klein-Gordon quantum equation from a completely metric origin. Thus, we have to look for possible – potentially approximated - solutions to the equation:

$$M^2 \cdot f \equiv \frac{m^2 \cdot c^2}{\hbar^2} \cdot f = \frac{R^*}{6} \cdot \left(f + \frac{C_f}{M^2} \right)^3. \quad (6)$$

This, however, is a polynomial of third order and it can have three solutions.

A bit of reshaping the last equation is going to help us to realize how this will possibly also solve the 3-generations mass problem:

$$\Psi \equiv M^2 \cdot f \equiv \frac{m^2 \cdot c^2}{\hbar^2} \cdot f = \frac{R^*}{6} \cdot \left(f + \frac{C_f}{M^2} \right)^3 = \frac{R^*}{M^6 \cdot 6} \cdot (\Psi + C_f)^3$$

$$\xrightarrow{R^{**} = \frac{R^*}{M^6 \cdot 6}; R^{***} = C_f} R^{**} \cdot (\Psi + R^{***})^3 - \Psi = 0 \quad (7)$$

$$= (\Psi + R^{***})^3 - \frac{\Psi}{R^{**}} = (R^{***})^3 - \frac{\Psi}{R^{**}} + 3 \cdot (R^{***})^2 \cdot \Psi + 3 \cdot R^{***} \cdot \Psi^2 + \Psi^3$$

The general solution to a three-order polynomial could be given via the following product form:

$$(\Psi - \Psi_1) \cdot (\Psi - \Psi_2) \cdot (\Psi - \Psi_3) =$$

$$\Psi^3 - \Psi^2 \cdot (\Psi_1 + \Psi_2 + \Psi_3) + \Psi \cdot (\Psi_1 \Psi_2 + \Psi_1 \Psi_3 + \Psi_2 \Psi_3) - \Psi_1 \Psi_2 \Psi_3 \cdot$$

$$(R^{***})^3 - \frac{\Psi}{R^{**}} + 3 \cdot (R^{***})^2 \cdot \Psi + 3 \cdot R^{***} \cdot \Psi^2 + \Psi^3 \quad (8)$$

Comparing the latter with the last line in (7) gives us:

$$3 \cdot R^{***} = 3 \cdot C_f = -(\Psi_1 + \Psi_2 + \Psi_3)$$

$$3 \cdot (R^{***})^2 - \frac{1}{R^{**}} = 3 \cdot (C_f)^2 - \frac{6 \cdot M^6}{R^*} = \Psi_1 \Psi_2 + \Psi_1 \Psi_3 + \Psi_2 \Psi_3 \cdot$$

$$(C_f)^3 = -\Psi_1 \Psi_2 \Psi_3 \quad (9)$$

Thus, we have obtained the three generations of quantum gravity solutions to the combined mass²-times-f-function, given via $\Psi \equiv M^2 \cdot f \equiv \frac{m^2 \cdot c^2}{\hbar^2} \cdot f$, as functions or dependencies of the Ricci scalar

R^* of the quantum-gravity (just scalar) variated metric $G_{\alpha\beta}=F[f]^*g_{\alpha\beta}$, the mass-values $M^2 \equiv \frac{m^2 \cdot c^2}{\hbar^2}$ and a constant C_f .

Please note: As the expression $\Psi_1\Psi_2 + \Psi_1\Psi_3 + \Psi_2\Psi_3$ on the right-hand side in the second line in (9) should be just a constant and thus, could not depend on the various masses M_i , we have to demand the Ricci-scalar R^* to be directly connected with the M_i via a constant const:

$$\frac{6 \cdot M^6}{R^*} = \text{const} \Rightarrow R^* = \frac{6 \cdot M^6}{\text{const}}. \quad (10)$$

This gives two equations for the extraction of two of the three states Ψ_i out of the other:

$$\begin{aligned} R^{***} = C_f &= -\frac{\Psi_1 + \Psi_2 + \Psi_3}{3} \\ \Rightarrow \left\{ \begin{array}{l} \frac{1}{3} \cdot (\Psi_1 + \Psi_2 + \Psi_3)^2 - \frac{6 \cdot M^6}{R^*} \\ = \frac{1}{3} \cdot (\Psi_1 + \Psi_2 + \Psi_3)^2 - \text{const} = \Psi_1\Psi_2 + \Psi_1\Psi_3 + \Psi_2\Psi_3 \end{array} \right. & \quad (11) \\ \Rightarrow \left(\frac{\Psi_1 + \Psi_2 + \Psi_3}{3} \right)^3 &= \Psi_1\Psi_2\Psi_3 \end{aligned}$$

We find with a vanishing constant const (the curvature term) we would obtain one solution, namely:

$$\Psi_1 = \Psi_2 = \Psi_3. \quad (12)$$

We also find that with a vanishing constant C_f (the shift of the function f with the value $F[f]=f=0$ moving away from the origin), we would only obtain the solution

$$\Psi_2 = \Psi_3 = \frac{1}{2} \left(-\Psi_1 \pm \sqrt{4 \cdot \text{const} - 3 \cdot \Psi_1^2} \right). \quad (13)$$

A bit more discussion shall be presented in [2].

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